

THE ZERO SET  
OF THE INDEPENDENCE POLYNOMIAL  
OF A GRAPH

MARTÍN SOMBRA (ICREA & UB)

Joint work with Juan Rivera-Letelier (Rochester)

COMPLEXITY OF NUMERICAL COMPUTATION

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**BGSMath**  
BARCELONA GRADUATE  
SCHOOL OF MATHEMATICS



**ICREA**



UNIVERSITAT DE  
BARCELONA

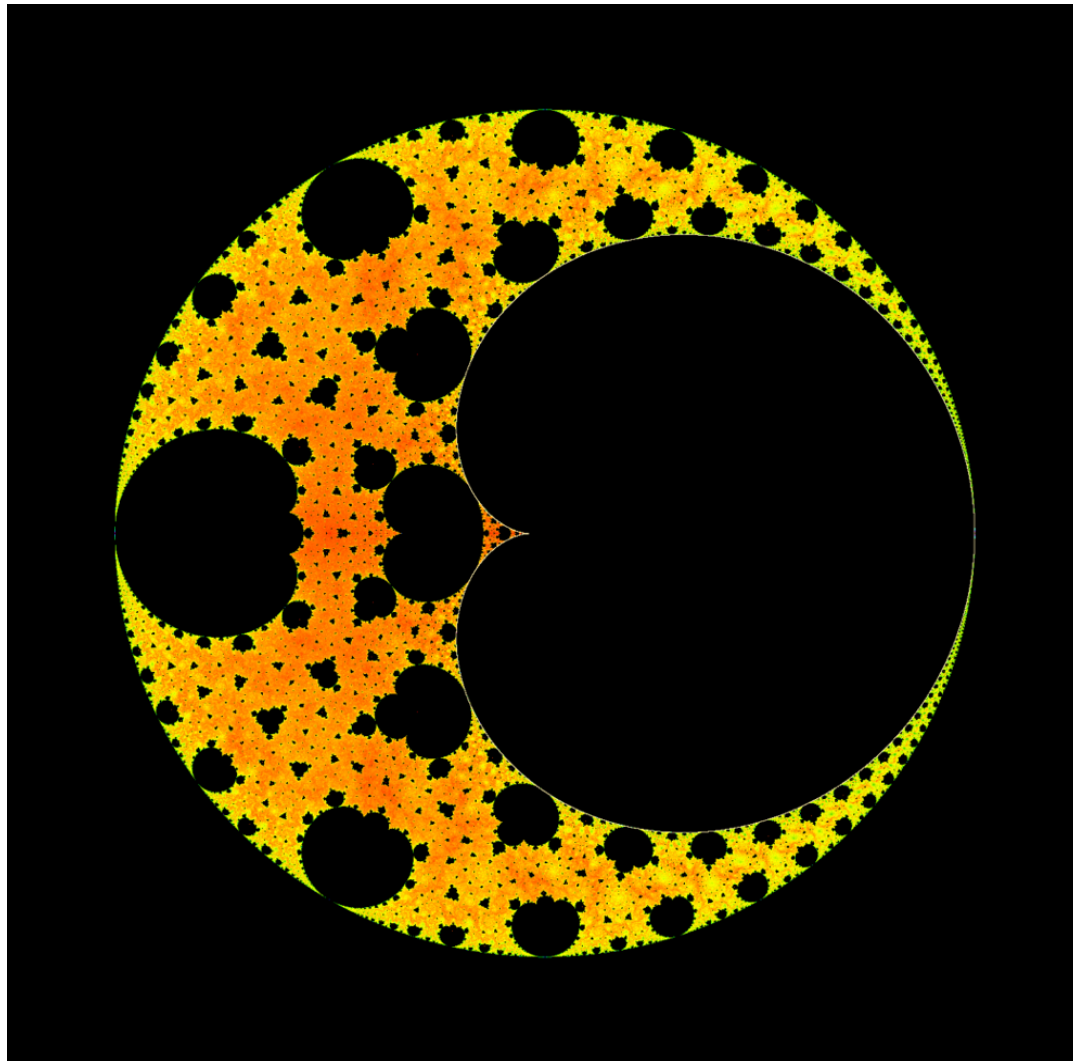


FIGURE BY N. FAGELLA

# INDEPENDENCE POLYNOMIALS OF GRAPHS

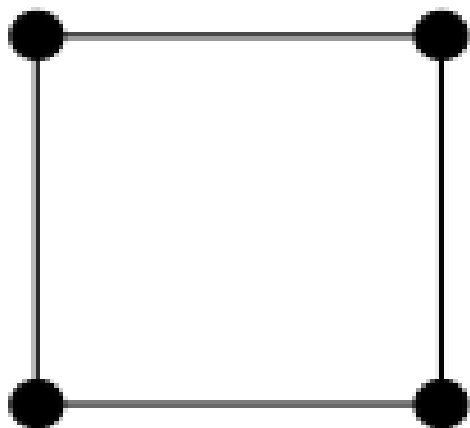
$G$  finite graph

$I \subset V_G$  independent set if  $v \sim v' \nRightarrow v, v' \in I$   
↑  
vertices of  $G$

The independence polynomial of  $G$  is

$$Z_G = \sum_{\substack{I \subset V_G \\ \text{indep}}} x^{\#I} \in \mathbb{Z}_{\geq 0}[x]$$

Ex

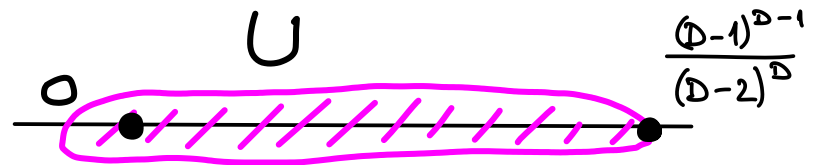


$$Z_G = 1 + 4x + 2x^2$$

$\S$  countably infinite graph of degree  $D \in [3, +\infty)$   
 $\uparrow$  mxl valence of nodes

THM (Peters & Regts 2019 after conjecture by Sokal 2001)

$\exists U \supset \left[0, \frac{(D-1)^{D-1}}{(D-2)^D}\right)$  open s.t.  $z_G(\lambda) \neq 0 \quad \forall G \in \mathcal{G}$  finite



Q: Is there a limit distribution for the zeros of  $z_G$ ?

More precisely, for  $F \in \mathbb{C}[x] \setminus \mathbb{C}$  set

$$\delta_F = \frac{1}{\deg(F)} \sum_{\lambda} \text{mult}_{\lambda}(F) \cdot \delta_{\lambda} \quad \text{discrete probab measure / } \mathbb{C}$$

$\exists \mu_{\mathcal{G}}$  probab measure /  $\mathbb{C}$  s.t.  $\lim_{G \rightarrow \mathcal{G}} \delta_{z_G} = \mu_{\mathcal{G}}$  weak- $\ast$ ?

# REGULAR ROOTED TREES

For  $d \geq 2$ :  $\mathbb{Z}$  regular rooted tree with branching  $d$

For  $k \geq -1$ :  $\mathbb{Z}_k$  subtree at depth  $\leq k$

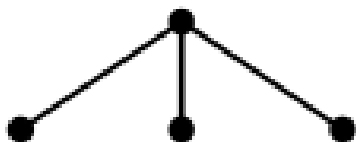
Ex:

$d=3$

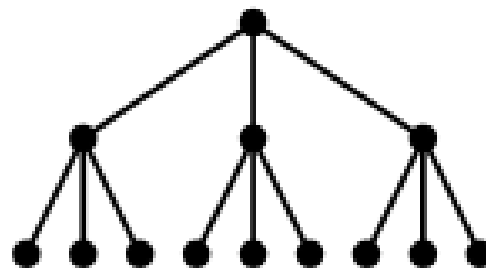
$k=0$



$k=1$



$k=2$



$T_k := \mathbb{Z}_{\mathbb{Z}_k}$  indep poly of  $\mathbb{Z}_k$

$$R_k := \frac{T_k}{T_{k-1}^d} \in \mathbb{Q}(x)$$

then  $T_{-1} = 1$ ,  $T_0 = 1+x$  &  $T_k = T_{k-1}^d + x \cdot T_{k-2}^d$  for  $k \geq 1$

and  $R_0 = 1+x$  &  $R_k = 1 + \frac{x}{R_{k-1}^d}$  for  $k \geq 1$

# COMPLEX DYNAMICS

For  $\lambda \in \mathbb{C}^*$

$$f_\lambda: \bar{\mathbb{C}} \rightarrow \bar{\mathbb{C}} \quad z \mapsto 1 + \frac{\lambda}{z^2}$$

Riemann sphere

- Critical points:  $\{0, \infty\}$
- $0 \rightarrow \infty \rightarrow 1 \rightarrow 1 + \lambda$
- $f_\lambda^{-1}(\infty) = \{0\}$

$$\underbrace{f \circ \dots \circ f}_k$$

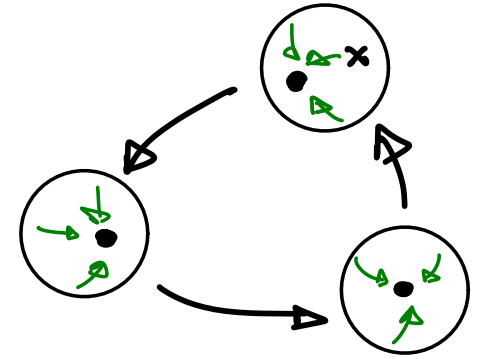
Hence  $R_k(\lambda) = f_\lambda^k(1 + \lambda) = f_\lambda^{k+3}(0)$

$$\Omega_k := (T_k = 0) = \{\lambda \in \mathbb{C}^* \mid f_\lambda^{k+3}(0) = 0\}$$

$$\Omega := \bigcup_{k \geq 0} \Omega_k = \{\lambda \in \mathbb{C}^* \mid 0 \in \text{Per}(f_\lambda)\}$$

# HYPERBOLIC COMPONENTS AND ZERO FREE REGIONS

$\lambda \in \mathbb{C}^*$  **hyperbolic** if the critical points of  $f_\lambda$  converge to attracting cycles



In our case:

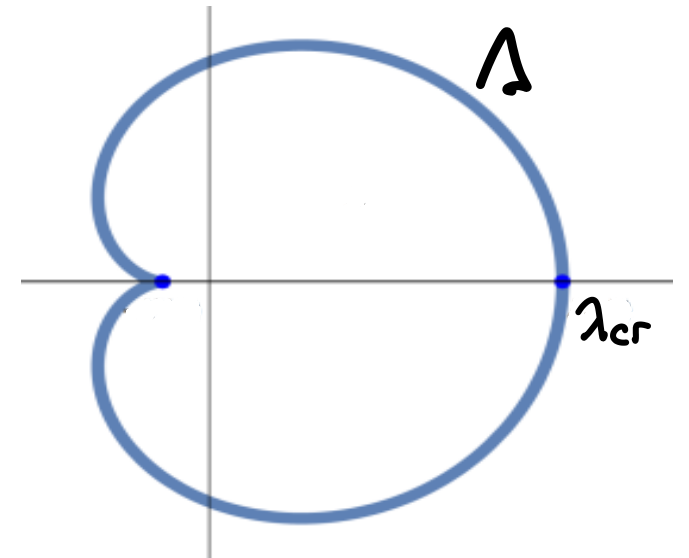
$\lambda$  hyperbolic iff  $f_\lambda$  has an attracting cycle (Fatou's theorem)

$\Rightarrow$  if  $\lambda$  is attracting but not superattracting then  $\lambda \notin \Omega$

$\Lambda := \{ \lambda \mid f_\lambda \text{ has } \underline{\text{attracting fixed point}} \}$

$$= \left\{ \frac{-d^d \alpha}{(d+\alpha)^{d+1}} \mid \alpha \in \mathbb{D} \right\}$$

$\Rightarrow \Lambda \cap \Omega = \emptyset$  (Peters & Regts 2019)



# A "BIFURCATION" HEIGHT

For  $\lambda \in \overline{\mathbb{Q}}^\times$  set

$$h(\lambda) := \lim_{k \rightarrow \infty} \frac{h_{\text{weil}}(f_\lambda^k(0))}{\deg(f_\lambda^k)} \in \mathbb{R}_{\geq 0} \quad (\text{if well-defn})$$

## THM (RL-S)

This limit defines a function  $h: \overline{\mathbb{Q}}^\times \rightarrow \mathbb{R}_{\geq 0}$  with

$$h(\lambda) = 0 \quad \text{iff} \quad 0 \in \text{PrePer}(f_\lambda)$$

Moreover  $h(\lambda) = \rho(\lambda) + \sum_{p \text{ prime}} \log^+ |\lambda|_p$  with  $\rho: \mathbb{C}^\times \rightarrow \mathbb{R}$

potential of  $\mu_{\text{bit}}$ , the normalized bifurcation measure

of  $(f_\lambda)_{\lambda \in \mathbb{C}^\times}$ , and  $\rho - \log^+$  is continuous on  $\overline{\mathbb{C}}$



# QUANTITATIVE EQUIDISTRIBUTION

$h$  is a (quasicanonical) height function (Arakelov geometry)  
+ arithmetic equidistribution thm (Favre & RL 2006)

Wasserstein distance

COR  $\exists c > 0$  st  $W(S_{T_k}, \mu_{b;f}) \leq c \left( \frac{\log k}{d^k} \right)^{1/2}$

I.e.  $\forall f: \mathbb{C} \rightarrow \mathbb{R}$  Lipschitz

$$\left| \frac{1}{\#\Omega_k} \sum_{\lambda \in \Omega_k} f(\lambda) - \int f d\mu_{b;f} \right| \leq c \left( \frac{\log k}{d^k} \right)^{1/2} \text{Lip}(f)$$

In particular  $\lim_{k \rightarrow \infty} S_{T_k} = \mu_{b;f}$

$$d = 2$$

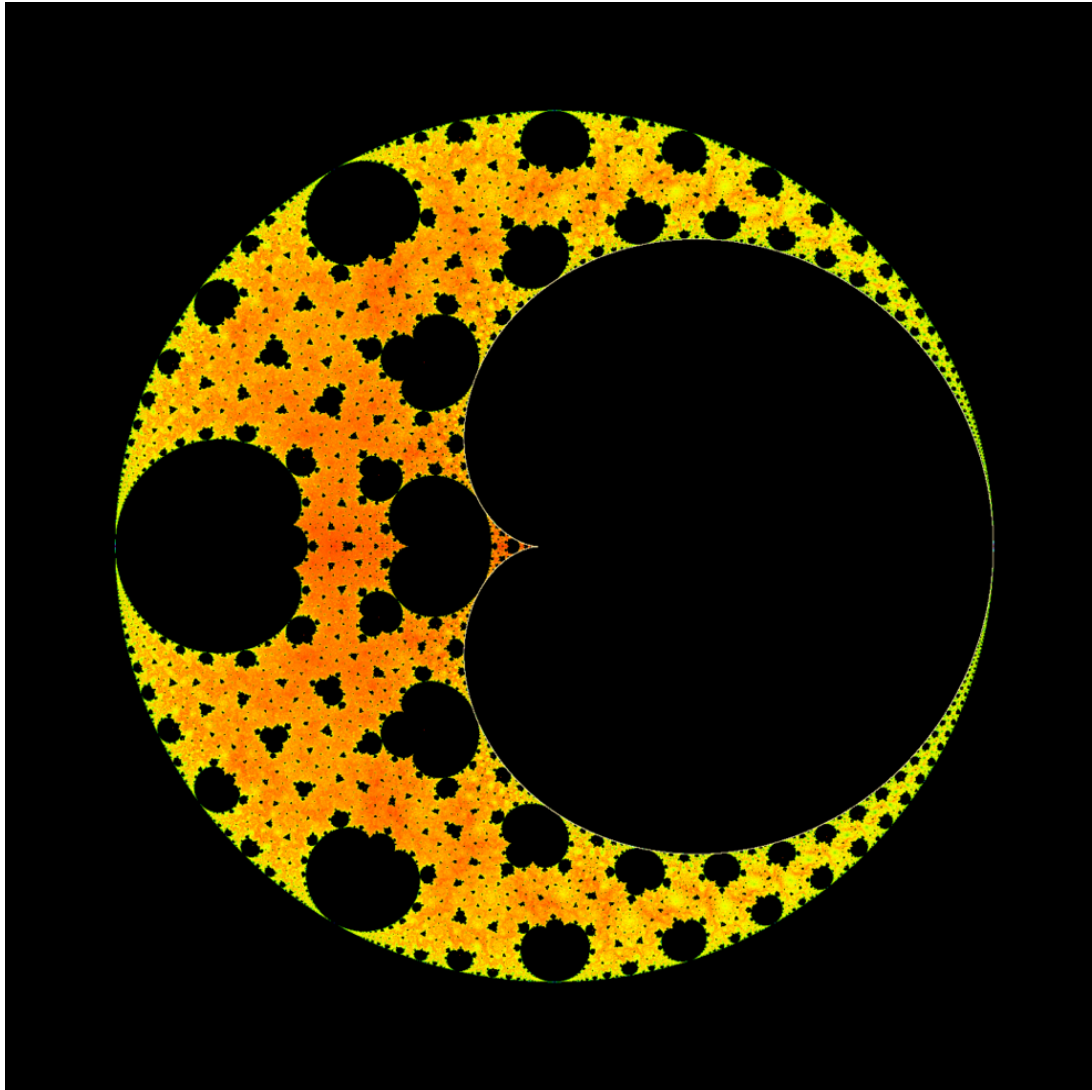


FIGURE BY N. FAGELLA

$$d = 3$$

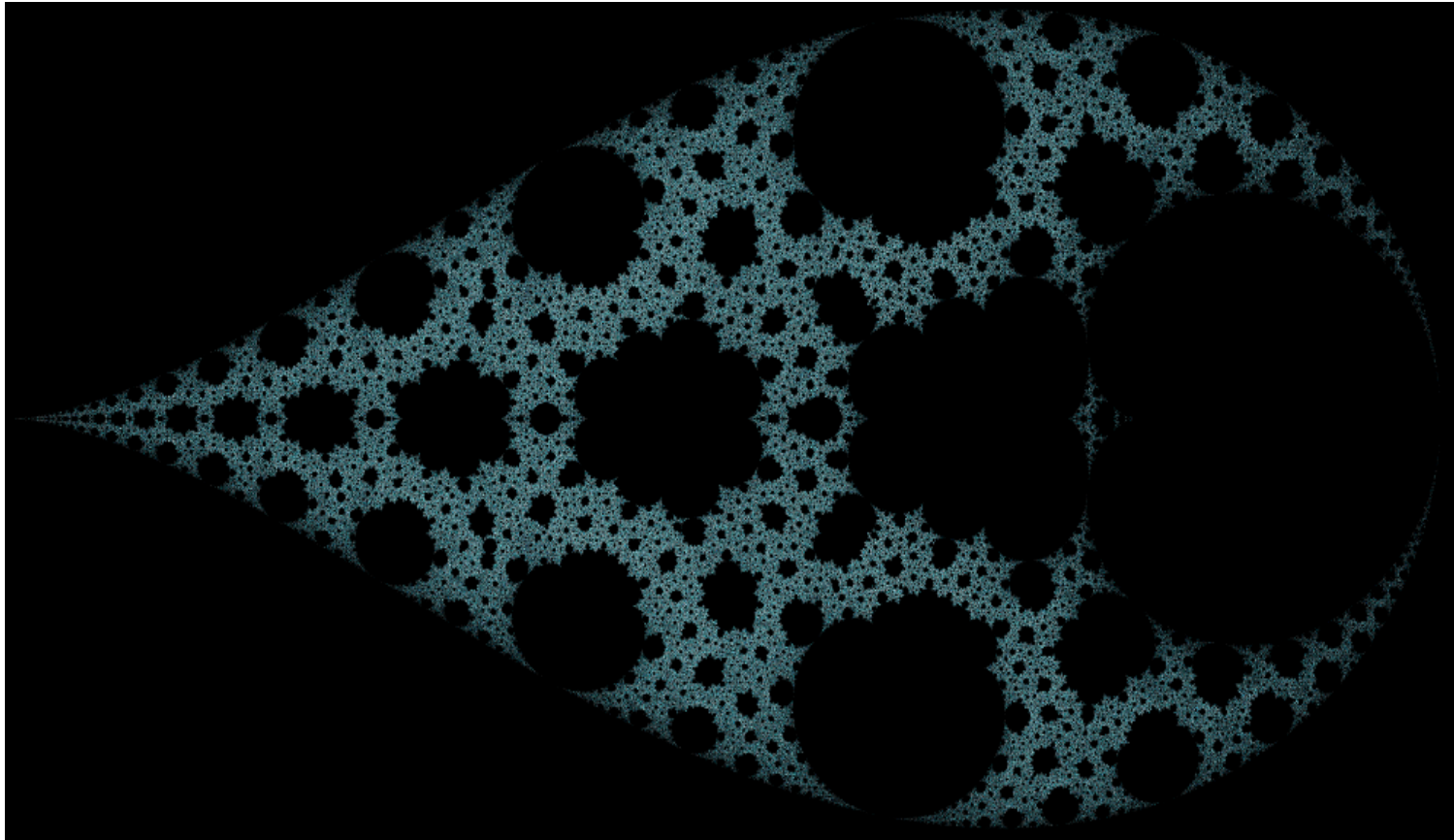
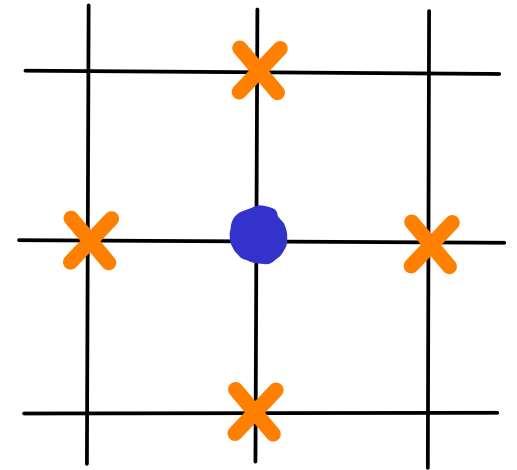


Figure by B. Espigule

# THE LATTICE GAS HARDCORE MODEL

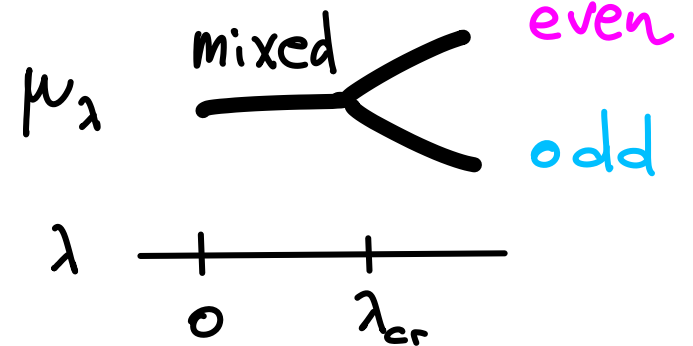
- Large particles on the vertices of  $\mathcal{G}$  excluding other particles at the adjacent ones
- Activity parameter  $\lambda \in [0, +\infty)$  determining the probability of particles appearing on vertices
- For  $G \subset \mathcal{G}$  finite, the partition function is  $Z_G$
- The pressure function is

$$P(\lambda) = \lim_{G \rightarrow \mathcal{G}} \frac{1}{\#V_G} \log(Z_G(\lambda))$$



# PHASE TRANSITIONS

Bifurcation of Gibbs measures



Related to *non-analyticity* of the pressure function

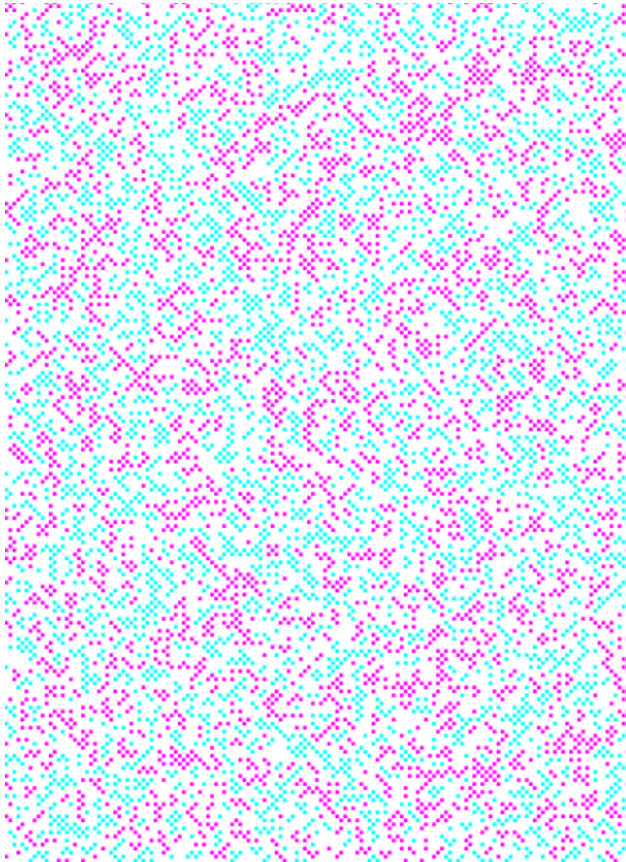
For the regular rooted tree  $\mathbb{Z}$ :

$$\lambda_{cr} = \frac{d^d}{(d-1)^{d+1}}$$

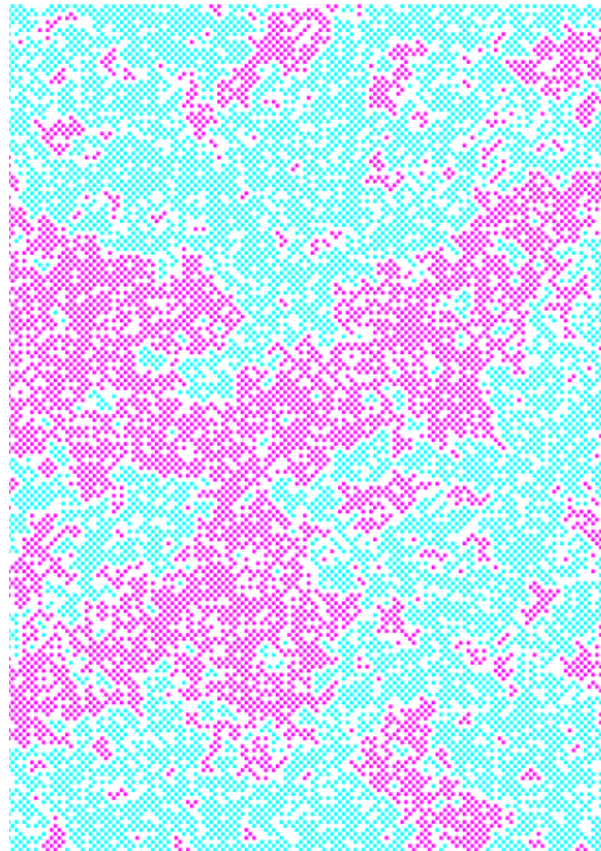
(Ruelle 1967)

For  $\mathbb{Z}^2$  :  $3.787 < \lambda_{cr} < 3.792$  (computationally!)

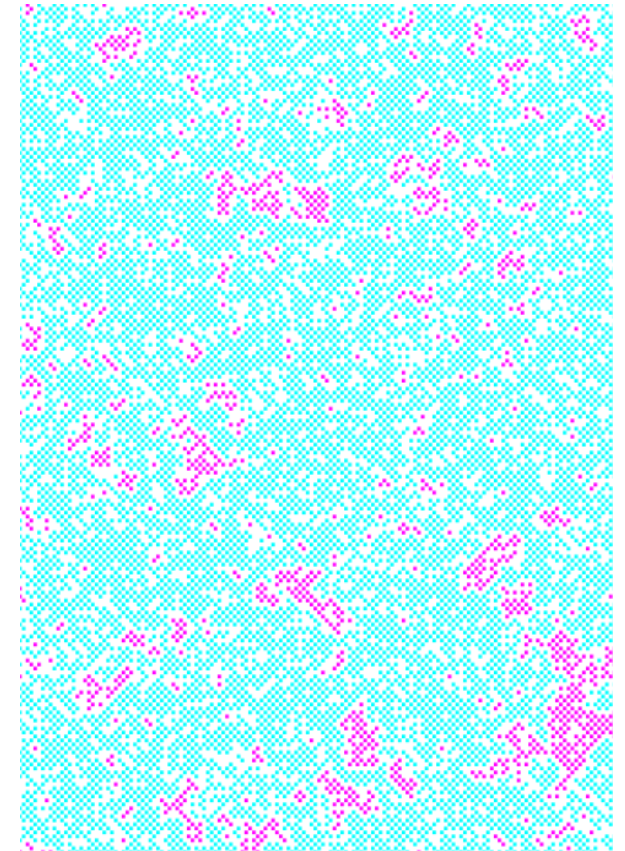
$\lambda = 1$



$\lambda = 3.787$



$\lambda = 3.792$



Figures by P. Shor & P. Winkler

# THE PRESSURE FUNCTION IS SMOOTH

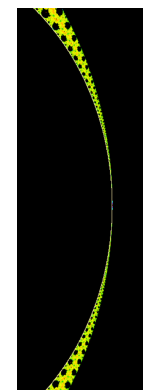
Set  $\mathcal{J} = \mathbb{C}$  again. The QET implies that

$$P(\lambda) = \int_{\mathbb{C}} \text{Log} |\lambda - z| d\mu_{\text{bit}}(z)$$

Used to prove:

PROP (RL-S)  $\exists c > 0$  s.t.

$$\forall \varepsilon > 0 \quad \mu_{\text{bit}}(B(\lambda_{cr}, \varepsilon)) \leq e^{-c/\varepsilon}$$



THM (RL-S)  $P$  is analytic on  $(0, +\infty) \setminus \{\lambda_{cr}\}$  and  $\mathcal{C}^\infty$  at  $\lambda_{cr}$

$\Rightarrow$  LGHM on  $\mathbb{C}$  has a phase transition of "infinite order" at  $\lambda_{cr}$

THANKS !